

**INTEGRAL MOMENTS  
AND BOUNDS FOR  
AUTOMORPHIC  
L-FUNCTIONS**

by

Adrian Diaconu

and

Paul Garrett

**1918:** Hardy–Littlewood

$$\int_0^T |\zeta(\tfrac{1}{2} + it)|^2 dt \sim T \log T$$

**1926:** Ingham

$$\int_0^T |\zeta(\tfrac{1}{2} + it)|^4 dt \sim \frac{1}{2\pi^2} T (\log T)^4$$

**1979:** Heath-Brown

$$\begin{aligned} \int_0^T |\zeta(\tfrac{1}{2} + it)|^4 dt \\ = T \cdot P_4(\log T) + \mathcal{O}\left(T^{\frac{7}{8} + \epsilon}\right) \end{aligned}$$

Let

$$f(z) = \sum_{n=1}^{\infty} a(n)e^{2\pi inz}$$

be a cusp form of weight  $\kappa$  for  $SL_2(\mathbb{Z})$  with associated  $L$ -function  $L_f(s)$ .

**1982:** Good

$$\begin{aligned} & \int_0^T \left| L_f \left( \frac{\kappa}{2} + it \right) \right|^2 dt \\ &= aT(\log T + b) + \mathcal{O} \left( (T \log T)^{\frac{2}{3}} \right) \end{aligned}$$

for certain constants  $a, b$ .

Good's result was extended to fourth moment of the Riemann zeta-function by Motohashi in 1993.

4

Let  $k$  be a number field

$$G = GL_2 \text{ over } k$$

and define the standard subgroups

$$P = \left\{ \begin{pmatrix} * & * \\ 0 & * \end{pmatrix} \right\}$$

$$N = \left\{ \begin{pmatrix} 1 & * \\ 0 & 1 \end{pmatrix} \right\}$$

$$M = \left\{ \begin{pmatrix} * & 0 \\ 0 & * \end{pmatrix} \right\}$$

$$Z = \text{center of } G$$

For any place  $\nu$  of  $k$ , let  $K_\nu^{\max}$  be the standard maximal compact subgroup; for finite  $\nu$ , we take

$$K_\nu^{\max} = GL_2(\mathfrak{o}_\nu)$$

at real places

$$K_\nu^{\max} = O(2)$$

at complex places

$$K_\nu^{\max} = U(2)$$

One of our goals is to form a suitable Poincaré series, and then unwind a corresponding global integral to express it as an inverse Mellin transform of an Euler product. The Poincaré series will be of the form

$$\text{Pé}(g) = \sum_{\gamma \in M_k \backslash G_k} \varphi(\gamma g) \quad (g \in G_{\mathbb{A}})$$

for suitable *monomial* vector

$$\varphi = \bigotimes_{\nu} \varphi_{\nu}$$

For *finite* primes  $\nu$ , the local component  $\varphi_\nu$  is defined as follows. For a character  $\chi_{0,\nu}$  of  $M_\nu$  ( $K_\nu = K_\nu^{\max}$ ), put

$$\varphi_\nu(g) = \begin{cases} \chi_{0,\nu}(m) & g = mk \in M_\nu \cdot K_\nu \\ 0 & g \notin M_\nu \cdot K_\nu \end{cases}$$

For  $\nu$  *infinite*, we need not entirely specify the local component  $\varphi_\nu$ , only requiring the left equivariance

$$\varphi_\nu(mnk) = \chi_{0,\nu}(m) \cdot \varphi_\nu(n)$$

for  $n \in N_\nu$  and  $m \in M_\nu$ . Require also that  $\chi_0 = \bigotimes_\nu \chi_{0,\nu}$  be  $M_k$ -invariant.

For the sake of clarity, we shall assume that

$$\chi_{0,\nu}(m) = \left| \frac{y_1}{y_2} \right|_\nu^v \quad m = \begin{pmatrix} y_1 & 0 \\ 0 & y_2 \end{pmatrix}, \quad v \in \mathbb{C}$$

Thus,  $\varphi$  has trivial central character. Then  $\varphi$  is left  $M_{\mathbb{A}}$ -equivariant by the character  $\chi_0$ . Also, for  $\nu$  infinite, our assumptions imply that the local component

$$x \rightarrow \varphi_\nu \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix}$$

is a function of  $|x|$  only.

Let  $f_1$  and  $f_2$  be cuspforms on  $G_{\mathbb{A}}$ .



The integral under consideration is

$$I(\chi_0) = \int_{Z_{\mathbb{A}} G_k \backslash G_{\mathbb{A}}} \text{Pé}(g) f_1(g) \bar{f}_2(g) dg$$

We assume *compatibility* among central characters so that the integrand is genuinely well-defined on  $Z_{\mathbb{A}} \backslash G_{\mathbb{A}}$ .

The integral immediately unwinds to

$$\int_{Z_{\mathbb{A}} M_k \backslash G_{\mathbb{A}}} \varphi(g) f_1(g) \bar{f}_2(g) dg$$

Using the Fourier expansion

$$f_1(g) = \sum_{\xi \in Z_k \backslash M_k} W_1(\xi g)$$

this further unwinds to

$$\int_{Z_{\mathbb{A}} \backslash G_{\mathbb{A}}} \varphi(g) W_1(g) \bar{f}_2(g) dg$$

Let  $C = GL_1(\mathbb{A})/GL_1(k)$  be the idele class group, and  $\widehat{C}$  its dual. More explicitly, by Fujisaki's Lemma,  $C$  is a product of a copy of  $\mathbb{R}^+$  and a compact group  $C_0$ . By Pontryagin duality,  $\widehat{C} \approx \mathbb{R} \times \widehat{C}_0$  with  $\widehat{C}_0$  discrete. Also, for any compact open subgroup  $U_{\text{fin}}$  of the finite-prime part in  $C_0$ , the dual of  $C_0/U_{\text{fin}}$  is finitely generated with rank  $[k : \mathbb{Q}] - 1$ .

With these notations, the general Mellin transform and inversion are

$$\begin{aligned}
 f(x) &= \int_{\widehat{C}} \int_C f(y) \chi(y) dy \chi^{-1}(x) d\chi \\
 &= \sum_{\chi' \in \widehat{C}_0} \frac{1}{2\pi i} \int_{\Re(s)=\sigma} \int_C f(y) \chi'(y) |y|^s dy \\
 &\quad \cdot \chi'^{-1}(x) |x|^{-s} ds
 \end{aligned}$$

for a suitably normalized Haar measure on  $C$ .

For  $\nu$  infinite and  $s \in \mathbb{C}$ ,

$$\begin{aligned} \mathcal{K}_\nu(s, \chi_{0,\nu}, \chi_\nu) = & \\ & \int_{Z_\nu \backslash M_\nu N_\nu} \int_{Z_\nu \backslash M_\nu} \varphi_\nu(m_\nu n_\nu) W_{1,\nu}(m_\nu n_\nu) \\ & \cdot \overline{W}_{2,\nu}(m'_\nu n_\nu) \chi_\nu(m'_\nu) |m'_\nu|_\nu^{s-\frac{1}{2}} \\ & \cdot \chi_\nu(m_\nu)^{-1} |m_\nu|_\nu^{\frac{1}{2}-s} dm'_\nu dn_\nu dm_\nu \end{aligned}$$

and set

$$\mathcal{K}_\infty(s, \chi_0, \chi) = \prod_{\nu|\infty} \mathcal{K}_\nu(s, \chi_{0,\nu}, \chi_\nu)$$

Here  $\chi_0 = \bigotimes_\nu \chi_{0,\nu}$  is the character defining  $\varphi$ , and  $\chi = \bigotimes_\nu \chi_\nu \in \widehat{C}_0$ .

**Theorem.** *For  $\varphi$  an admissible monomial vector, for suitable  $\sigma > 0$ , we have*

$$I(\chi_0) = \sum_{\chi \in \widehat{C}_0} \frac{1}{2\pi i} \int_{\Re(s)=\sigma} L(\chi_0 \cdot \chi^{-1} |\cdot|^{1-s}, f_1) \cdot L(\chi |\cdot|^s, \bar{f}_2) \mathcal{K}_\infty(s, \chi_0, \chi) ds$$

*Let  $S$  be a finite set of places including the archimedean places, all the absolutely ramified primes, and all the finite bad places for  $f_1$  and  $f_2$ . Then, in fact, the sum is over a set  $\widehat{C}_{0,S}$  of characters unramified outside  $S$  with bounded ramification at finite places.*

Simplifying assumption: Both  $f_1$  and  $f_2$  have trivial central character.

**Proof (sketch):** Applying the Mellin transform and inversion to  $\bar{f}_2$  via the obvious identification  $Z_{\mathbb{A}} M_k \backslash M_{\mathbb{A}} \approx C$ , and using the Fourier expansion

$$f_2(g) = \sum_{\xi \in Z_k \backslash M_k} W_2(\xi g)$$

the integral  $I(\chi_0)$  is

$$\int_{Z_{\mathbb{A}} \backslash G_{\mathbb{A}}} \varphi(g) W_1(g) \cdot \left( \int_{\widehat{C}} \int_{Z_{\mathbb{A}} M_k \backslash M_{\mathbb{A}}} \bar{f}_2(m'g) \chi(m') dm' d\chi \right) dg$$

$$\begin{aligned}
&= \int_{\widehat{C}} \left( \int_{Z_{\mathbb{A}} \backslash G_{\mathbb{A}}} \varphi(g) W_1(g) \cdot \right. \\
&\quad \left. \cdot \int_{Z_{\mathbb{A}} \backslash M_{\mathbb{A}}} \overline{W}_2(m'g) \chi(m') dm' dg \right) d\chi
\end{aligned}$$

We remark that all the characters  $\chi$  which appear are unram. outside  $S$  with bounded ramification.

Assuming at this point that  $f_1$  and  $f_2$  generate irreducible representations locally everywhere, the Whittaker functions  $W_i$

factor over primes

$$W_i(\Pi_\nu g_\nu) = \Pi_\nu W_{i,\nu}(g_\nu)$$

Therefore, the inner integral over  $Z_{\mathbb{A}} \backslash G_{\mathbb{A}}$  and  $Z_{\mathbb{A}} \backslash M_{\mathbb{A}}$  factors over primes, and

$$I(\chi_0) = \int_{\widehat{C}} \Pi_\nu \left( \int_{Z_\nu \backslash G_\nu} \int_{Z_\nu \backslash M_\nu} \varphi_\nu(g_\nu) \cdot W_{1,\nu}(g_\nu) \overline{W}_{2,\nu}(m'_\nu g_\nu) \chi_\nu(m'_\nu) dm'_\nu dg_\nu \right) d\chi$$

Assume that  $\nu \notin S$ . Then,  $f_1$  and  $f_2$  are spherical at  $\nu$ . It can be shown that the



local integral

$$= L_\nu(\chi_{0,\nu} \cdot \chi_\nu^{-1} |\cdot|_\nu^{1/2}, f_1) \cdot L_\nu(\chi_\nu |\cdot|_\nu^{1/2}, \bar{f}_2)$$

i.e., the product of local factors of the standard  $L$ -functions in the theorem.

**Remark.** The archimedean local factors are not at all specified. The option to vary the choices is essential.

**Theorem.** *The Poincaré series  $Pé(g)$  has meromorphic continuation in  $\chi_0$  and in suitable archimedean data.*

The idea of the proof is that after an obvious Eisenstein series is subtracted, the

Poincaré series is in  $L^2$  and is of sufficiently rapid decay in Siegel sets so that we can write its  $L^2$  spectral decomposition in terms of cuspforms and continuous spectrum.

For instance, let  $f$  be a cuspform on  $G_{\mathbb{A}}$  which everywhere is a spherical Hecke eigenfunction. The integral

$$\int_{Z_{\mathbb{A}} G_k \backslash G_{\mathbb{A}}} \bar{f}(g) P_{\nu}(g) dg$$

is an Euler product. At a finite  $\nu$ , the corresp. local factor is  $L_{\nu}(\chi_{0,\nu} | \cdot |_{\nu}^{1/2}, \bar{f})$  (up to a constant depending on the set of absolutely ramified primes in  $k$ ).

## Applications

For  $\nu, w \in \mathbb{C}$ , define

$$\varphi_\nu \left( \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \right) = (x^2 + 1)^{-\frac{w}{2}}$$

for  $\nu | \infty$  real,

$$\varphi_\nu \left( \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \right) = (|x|^2 + 1)^{-w}$$

for  $\nu | \infty$  complex, and choose

$$\chi_{0,\nu}(m) = |y|_\nu^v \quad m = \begin{pmatrix} y & 0 \\ 0 & 1 \end{pmatrix} \in M_\nu$$

The corresponding monomial vector  $\varphi$  is admissible for  $\Re(\nu)$  and  $\Re(w)$  sufficiently large.

The integral  $I = I(\chi_0) = I(v, w)$  can be written as

$$I = \sum_{\chi \in \widehat{C}_{0,S}} \frac{1}{2\pi i} \int_{\Re(s)=\sigma} L(\chi^{-1} | \cdot |^{v+1-s}, f_1) \cdot L(\chi | \cdot |^s, \bar{f}_2) \mathcal{K}_\infty(s, v, w, \chi) ds$$

Recall that  $\mathcal{K}_\infty = \prod_{\nu|\infty} \mathcal{K}_\nu$ . The function  $\mathcal{K}_\nu(s, v, w, \chi_\nu)$  is analytic in a region  $\mathcal{D} : \Re(s) = \sigma > \frac{1}{2} - \epsilon_0$ ,  $\Re(v) > -\epsilon_0$  and  $\Re(w) > \frac{3}{4}$ , with a fixed (small)  $\epsilon_0 > 0$ .

Moreover,

$$\begin{aligned} \mathcal{K}_\nu(s, v, w, \chi_\nu) &= A(v, w, \mu_1, \mu_2) \cdot \\ &\quad \cdot (1 + \ell_\nu^2 + 4(t + t_\nu)^2)^{-w} \cdot \\ &\quad \cdot \left[ 1 + \mathcal{O} \left( \left( \sqrt{1 + \ell_\nu^2 + 4(t + t_\nu)^2} \right)^{-1} \right) \right] \end{aligned}$$

for  $\nu|\infty$  complex, and

$$\begin{aligned} \mathcal{K}_\nu(s, v, w, \chi_\nu) &= B(v, w, \mu_1, \mu_2) \cdot \\ &\quad \cdot (1 + |t + t_\nu|)^{-w} \cdot \\ &\quad \cdot \left[ 1 + \mathcal{O} \left( (1 + |t + t_\nu|)^{-\frac{1}{2}} \right) \right] \end{aligned}$$

for  $\nu|\infty$  real. The functions  $A$  and  $B$  are ratios of gamma functions.

It follows that for  $\Re(w)$  sufficiently large,

$$I(0, w) = \sum_{\chi \in \widehat{C}_{0,S}} \frac{1}{2\pi} \int_{-\infty}^{\infty} L(\chi^{-1} | \cdot |^{\frac{1}{2}-it}, f_1) \\ \cdot L(\chi | \cdot |^{\frac{1}{2}+it}, \bar{f}_2) \mathcal{K}_{\infty}(\frac{1}{2} + it, 0, w, \chi) dt$$

**Remark.** Since  $I(0, w)$  has meromorphic continuation, a mean value result can be established by standard arguments.

**Remark.** When  $k = \mathbb{Q}$ , one recovers the classical results.

**Example.** Assume  $k = \mathbb{Q}(i)$ . For  $\ell \in \mathbb{Z}$ , define  $\Psi_{\infty}^{4\ell}(m) = (m/|m|)^{4\ell}$ .

As  $|T| \rightarrow \infty$ , we have the asymptotic formula

$$\sum_{|\ell| < \frac{T}{2}} \int_0^{\sqrt{T^2 - 4\ell^2}} |L\left(\frac{1}{2} + it, f \otimes \Psi_\infty^{4\ell}\right)|^2 dt \sim C T^2 \log T$$

where  $C$  is a non-zero (computable) constant (joint with Dorian Goldfeld).

**Corollary (Sarnak 1985).** *We have the following estimate:*

$$\sum_{|\ell| < T} \int_0^T |L\left(\frac{1}{2} + it, f \otimes \Psi_\infty^{4\ell}\right)|^2 dt \ll T^2 \log T$$

**Remark.** By making a careful choice of the monomial vector  $\varphi$ , one can exhibit an error term on the order of  $\mathcal{O}(T^\vartheta)$ , with  $\vartheta < 2$ . in the above asymptotic formula. One can deduce that

$$\left| L\left(\frac{1}{2} + it, f \otimes \Psi_\infty^{4\ell}\right) \right| \ll_\epsilon |t|^{\frac{\vartheta}{2} + \epsilon}$$

$$\left| L\left(\frac{1}{2} + it, f \otimes \Psi_\infty^{4\ell}\right) \right| \ll_\epsilon |\ell|^{\frac{\vartheta}{2} + \epsilon}$$



In the general case, for  $\chi \in \widehat{C}_0$  and  $t \in \mathbb{R}$ , put

$$\begin{aligned} \kappa_\chi(t, w) &= \prod_{\substack{\nu|\infty \\ \nu \text{ real}}} (1 + |t + t_\nu|)^{-w} \cdot \\ &\cdot \prod_{\substack{\nu|\infty \\ \nu \text{ complex}}} \left( \sqrt{1 + \ell_\nu^2 + 4(t + t_\nu)^2} \right)^{-w} \end{aligned}$$

where  $t_\nu$  and  $\ell_\nu$  are the parameters of the local component  $\chi_\nu$  of  $\chi$ . Since  $\chi$  is trivial on positive reals, we must have

$$\sum_{\nu|\infty} \alpha_\nu t_\nu = 0 \quad (\alpha_\nu = 1, 2)$$

The asymptotic formula one obtains is of type

$$\sum_{\chi \in \widehat{C}_{0,S}} \int_{\mathfrak{I}_\chi(T)} |L(\tfrac{1}{2} + it, f \otimes \chi)|^2 dt \sim C T^{[k:\mathbb{Q}]} (\log T)^{r_1+r_2}$$

for a non-zero positive constant  $C$ , where

$$\mathfrak{I}_\chi(T) = \{t : 0 \leq \kappa_\chi(t, -1) \leq T^{r_1+r_2}\}$$